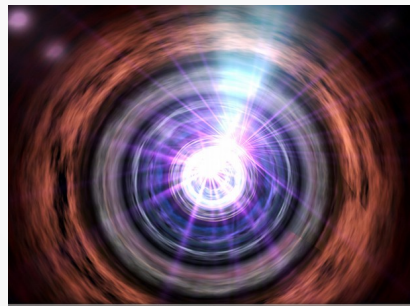


Is the Blazar Sequence related to accretion disk winds?



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The Model

The kinetic equation for electrons is described by

$$\frac{\partial n_e(\gamma, t)}{\partial t} + \frac{n_e(\gamma, t)}{t_{esc}} = Q_{e, inj} + L_e$$

where n_e is the differential electron density and Q_e , L_e are the injection and energy loss operations for electrons respectively. These are defined as

$$Q_{e, inj} = \begin{cases} k_{e1} \gamma^{-p} & \text{for } \gamma_{min} \leq \gamma \leq \gamma_{br} \\ k_{e2} \gamma^{-q} e^{-\gamma/\gamma_{max}} & \text{for } \gamma_{br} \leq \gamma \leq \gamma_{max}, \end{cases}$$

$$\gamma_{max} = \frac{3m_e c^2}{4\sigma_\tau c t_{acc} U_{tot}},$$

$$L_e = \frac{4}{3} \frac{\sigma_\tau}{m_e c^2} \frac{\partial}{\partial \gamma} [\gamma^2 n_e(\gamma, t) U_{tot}],$$

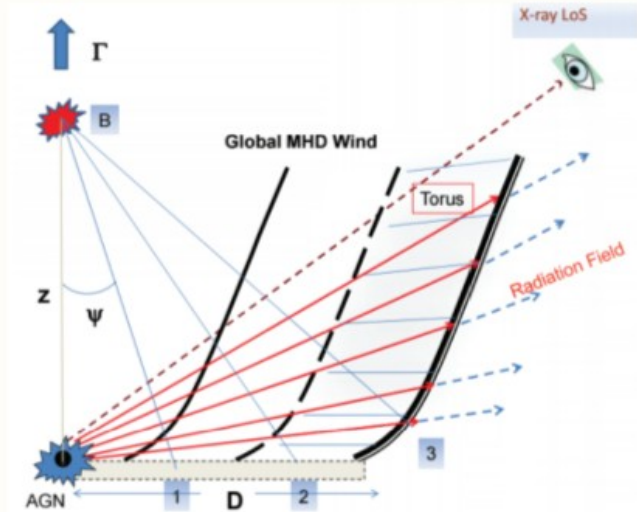
where $\gamma_{br} = 0.1\gamma_{max}$ and U_{tot} the total energy density. This is given by

$$U_{tot} = U_B + U_{ext} + U_{ssc}$$

The acceleration time t_{acc} is defined from the first order Fermi acceleration as

$$t_{acc} = \tau_{FI} \geq \frac{6\gamma m_e c^3}{e B u_s^2}$$

where u_s the velocity of the shock.



The wind's particle density is

$$n = n_0 \frac{r_0}{r}$$

where $r_0 = r_s$, $n_0 = \frac{\eta_w \dot{m}}{2\sigma_\tau r_s}$, $\dot{m} = \frac{\dot{M}}{\dot{M}_{Edd}}$ the mass rate normalized to Eddington mass rate, r_s is Schwarzschild radius, η_w the ratio of the mass-outflow rate in the wind to the mass-accretion rate \dot{m} , assumed here to be $\eta_w \simeq 1$ and σ_τ the Thomson cross section.

Photons from the accretion disk could be scattered on wind particles. We assume a spherical region between radii R_1 and R_2 , therefore the optical depth in first order is

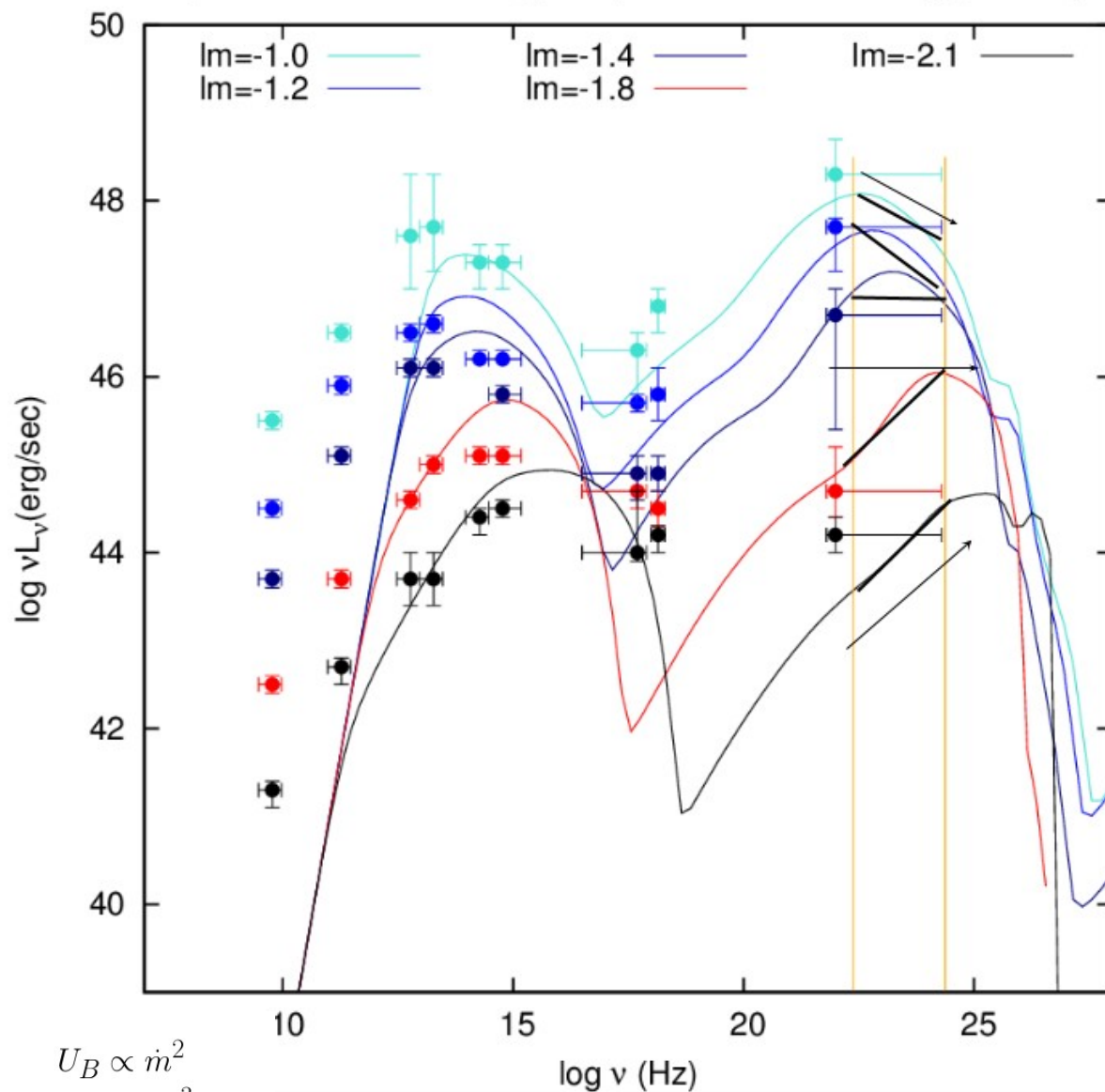
$$\tau_\tau = \int_{R_1}^{R_2} n(r) \sigma_\tau dr = n_0 \sigma_\tau \ln \frac{R_2}{R_1}$$

In the case of $\tau_\tau \ll 1$ the photon energy density of the scattered accretion disk photons up to distance R_2 is

$$U_{ext} = \frac{\eta_{ext} L_{acc} \tau_\tau}{4\pi R_2^2 c},$$

where η_{ext} the radiation to accretion efficiency, $L = \epsilon \dot{m}^2 L_\odot \hat{M}$ the ionizing luminosity from the disk and $L_\odot = 1.3 \times 10^{38} \text{ erg s}^{-1}$ is the Eddington luminosity, ϵ is the efficiency of conversion of mass into radiation for $\dot{m} = 1$ and $\hat{M} = \frac{\dot{M}}{M_\odot}$, where M_\odot is one solar mass.

Spectrum for different $\log(\dot{m}) = l_m$, $R = 9 \times 10^{15} \text{ cm}$, $M_{\text{BH}} = 10^9 M_\odot$



$$U_B \propto \dot{m}^2$$

$$U_{\text{ext}} \propto \dot{m}^3$$

$$\gamma_{\text{max}} \propto \dot{m}^{-2}(1 + \dot{m})$$

$$L_e^{\text{inj}} \propto \dot{m}^2$$

$$t_{\text{acc}} \propto \dot{m}^0$$

$\log \dot{m}$	B (G)	$\log \gamma_{br}$	$\log l_{\text{ext}}$	$\log l_e$
-1.0	2.40	2.90	-2.0	-1.9
-1.2	1.51	3.01	-2.1	-2.3
-1.4	0.95	3.16	-2.5	-2.7
-1.8	0.38	3.50	-3.4	-3.5
-2.1	0.19	3.60	-4.5	-4.1

$1 - \alpha_\gamma$	α_γ	Γ_γ	$\log(L_\gamma)$ (erg/sec)
-0.28	1.28	2.28	48.1
-0.30	1.30	2.30	47.7
-0.05	1.05	2.05	47.2
0.24	0.76	1.76	46.0
0.45	0.55	1.55	44.7

